# Fully Implicit Filtered P<sub>N</sub> for High-Energy



Vincent M. Laboure, Ryan G. McClarren, Cory D. Hauck

Department of Nuclear Engineering, Texas A&M University, College Station, TX, USA 77843

### **Motivation**

- ► All deterministic methods to solve the transport equation have flaws
- ► The PN equations produce a rotationally-invariant solution and therefore is **immune to ray effects** (good for diffusive problems)
- ► BUT they can also create **unphysical oscillations** and **negativity** in the solution
- ► While it is often crucial to have positivity in numerous applications (radiative transfer,...)
- ► Idea: **damp** the large derivatives in angle

## **Fully-Implicit** $P_N$ Equations

### ► Transport equation:

### **Eigenspectrum of** L + C

### Properties of the filters:

- ► All the eigenvalues are such that  $\Re\lambda\geq 1$
- Smallest eigenvalue  $\equiv \lambda_{\min} = 1$
- Largest eigenvalue  $\equiv \lambda_1$





## Impact on the eigenspectrum in void

 $c\Delta t$ В  $\lambda_{\min}$ 

 $\frac{1}{c}\partial_t\psi + \mathbf{\Omega}\cdot\nabla\psi + \sigma_t\psi = \sigma_s\phi + Q$ 

Spherical Harmonics expansion:

 $\psi(\mathbf{r}, \mathbf{\Omega}) \approx \sum_{l=1}^{N} \sum_{l=1}^{l} \psi_{l}^{m}(\mathbf{r}) Y_{l}^{m}(\mathbf{\Omega})$ l=0 m=-1

- Quasi steady-state formulation of the transport equation:  $A_{x}\partial_{x}\psi^{n+1} + A_{y}\partial_{y}\psi^{n+1} + A_{z}\partial_{z}\psi^{n+1} + \sigma_{t}^{*}\psi^{n+1} = \sigma_{s}\phi^{n+1} + Q^{n,*}$
- ► Linear system solved using a Krylov solver (PETSc):  $L\psi + C\psi = Q^*$

#### Filtering: McClarren/Hauck approach

► After each time step:

$$\psi_l^m \longleftarrow \frac{\psi_l^m}{1 + \alpha l^2 (l+1)^2}$$

Where:

$$\alpha = \frac{\omega}{N^2(\sigma_t L + N)^2}$$

$\Delta X$				
1	0	8.0868	1.0000	100%
1	20	8.0842	1.0840	91.1%
1	60	8.0788	1.1591	84.2%
1	100	8.0727	1.0981	89.6%
10	0	71.868	1.0000	100%
10	20	71.842	1.8404	53.7%
10	60	71.788	2.5911	37.7%
10	100	71.727	1.9807	49.7%

# $\gamma \equiv \frac{|\lambda_1 - \lambda_{\min}|}{|\lambda_{\min}|}$

 $\Gamma \equiv - \gamma$  $\gamma$ unfiltered

Phi 1.58e+018 <mark>\_</mark>\_

-1.58e+015 <sup>=</sup>0

\_8e+17

### **Thermal Radiation Transfer**

► Add temperature equation:

$$\begin{cases} \frac{1}{c} \partial_t I + \mathbf{\Omega} \cdot \nabla I + \sigma_t I = \sigma_a B(T) + \sigma_s \varphi + Q \\ C \partial_t T = \sigma \left( (\alpha - A \pi B(T)) + O \right) \end{cases}$$

 $\omega$ : filter strength L: characteristic length

#### Filtering: Radice/Abdikamalov approach

► Matrix form:

 $L\psi + (C + \Delta)\psi = Q^*$ 

Where:

$$\Delta \equiv -\beta \log \sigma_{\text{filterType}} \left(\frac{1}{N+1}\right)$$

And:

$$\sigma_{\text{filterType}} = \begin{cases} \sigma_{\text{Lanczos}}(\eta) = \frac{\sin \eta}{\eta} \\ \\ \sigma_{\text{SSpline}}(\eta) = \frac{1}{1 + \eta^4} \end{cases}$$

### **Effects of the filtering**

► Pulsed line source after  $10\Delta t$  (c = 1,  $\sigma_t = \sigma_s = 1$ ,  $\Delta t = 0.1$ , P1)

 $( C_v O_t I = O_a (\varphi - 4\pi D(I)) + Q_e$ 

### **Crooked Pipe Test Problem**

► Scalar flux after	$1000\Delta t~(\Delta t=4 imes 10^{-12}$ s, P	1)
$\beta =$	- 0	

 $\beta = 50$ 



 $\beta = 100$ 



 $\beta = 500$ 





 $\omega = 0.2$ 

 $\omega = 0$ 

Properties of the filters:

- ► do not change the 0th moment (particle conservation)
- ▶ vanish as  $N \to \infty$
- ► preserve the equilibrium diffusion limit
- ► preserve the rotational invariance



- Helps smooth the solution
- ► The strength of the filter should not be too large

### Conclusions

Filtered PN reduces the negativity and oscillations of the solution ► Filtered PN **improves the conditioning** of the linear system  $c\Delta t$ • Optimum filter strength in void depending on  $\frac{C\Delta t}{\Lambda}$  $\Delta x$ ► Reduces the speed of the waves

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